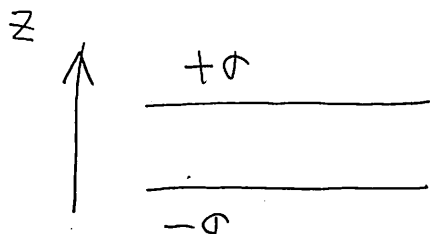


0.

(a) コレクターの2つの極板との表面電荷密度は $\pm\sigma$ ($\sigma = Q/S$)



左図のようにz軸をとる*

$$D = D \hat{z} \quad D = \sigma$$

電場は

$$E_1 = E_1 \hat{z} \quad E_1 = D/\epsilon_1 = \sigma/\epsilon_1$$

$$E_2 = E_2 \hat{z} \quad E_2 = D/\epsilon_2 = \sigma/\epsilon_2$$

(b) 電位差 V は

$$V = E_1 d_1 + E_2 d_2 = \sigma \left(\frac{d_1}{\epsilon_1} + \frac{d_2}{\epsilon_2} \right)$$

$$= \frac{Q}{S} \left(\frac{d_1}{\epsilon_1} + \frac{d_2}{\epsilon_2} \right)$$

* 対称性が2の形を仮定し
(端の効果は無視)

$\nabla \cdot D = \sigma \delta(z)$ と極板付近に体積積分すればわかる。

1.

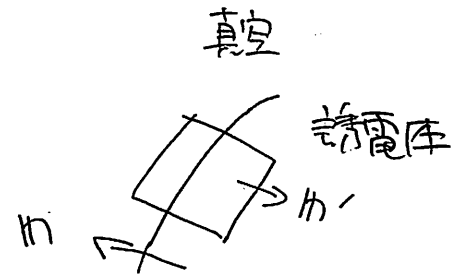
(a) 電荷を含む微小領域 dV の

$$\int_{dV} dV \cdot \rho_p(r) = - \int_{dV} dV \nabla \cdot P = - \int_{dS} dS \hat{n} \cdot P$$

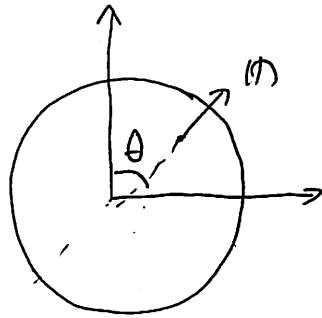
$$\hat{n}' = -\hat{n}$$

従って

$$\sigma_p(r) = \hat{n}(r) \cdot P(r)$$



(b)



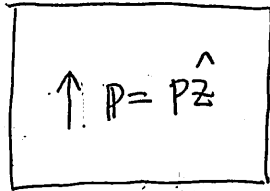
球座標系を用いる

$$\sigma_p = \hat{n} \cdot P = P \cos \theta$$

逆向型に一樣に分極した球

無極に似たように分極した系

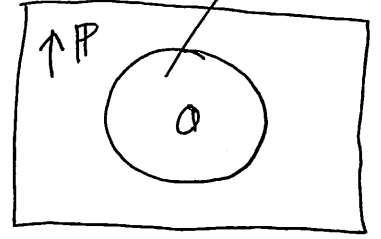
$$P = -P\hat{z}$$



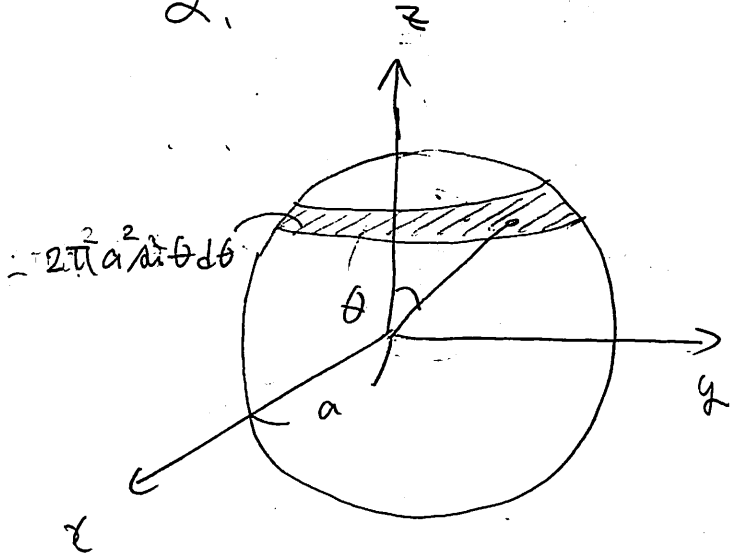
+



=



2.



表面上の点 $(a \cos\theta, a \sin\theta \cos\varphi, a \sin\theta \sin\varphi)$

その表面分極電荷は $\sigma_p = \mathbf{n} \cdot (-P)\hat{z} = -P \cos\theta$ ($\hat{n} = \hat{r}$)

これが穴の中に作る電場は

$$\begin{aligned} \mathbf{E}(0) &= \frac{1}{4\pi\epsilon_0} \int_0^\pi d\theta \, 2\pi a^2 \sin\theta \frac{-P \cos\theta}{a^2} (-\hat{r}) \\ &= \frac{P}{2\epsilon_0} \int_0^\pi d\theta \, \cos^2\theta \sin\theta \hat{z} = \frac{P}{3\epsilon_0} \hat{z} \end{aligned}$$

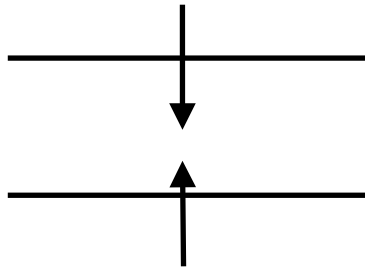
(1) $E = E_0 \hat{z}$ $E_0 = \sigma_0 / \epsilon_0$ 電位差は $\Delta\phi = E_0 l_z = \frac{\sigma_0}{\epsilon_0} l_z = \frac{Q}{\epsilon_0} \frac{l_z}{lxly}$
 かつ 静電容量は $C = Q / \Delta\phi = \epsilon_0 \frac{lxly}{l_z}$

静電場のエネルギーは

$$U_0 = \frac{1}{2} \epsilon_0 E_0^2 lxlyl_z =$$

$$F_z = - \frac{\partial U_0}{\partial l_z} = - \frac{1}{2} \epsilon_0 E_0^2 lxly$$

単位面積あたりに垂直に $T_3 = \frac{F_z}{lxly} = \frac{\epsilon_0}{2} E_0^2$
 圧力は図の向き。



(2)

$\phi_1 = \phi_2$ (境界面に垂直なので) $\vec{E}_1 = E_1 \hat{z}$ $\vec{E}_2 = E_2 \hat{z}$ $E_1 = \sigma/\epsilon_0 = E_0$
 $E_2 = \sigma/\epsilon = E_0 \frac{\epsilon_0}{\epsilon}$

電位差は $\Delta\phi = E_1(l_z - z) + E_2 z$

$= 0 \left[\frac{l_z}{\epsilon_0} + \left(\frac{1}{\epsilon} - \frac{1}{\epsilon_0} \right) z \right] = 0 \frac{l_z}{\epsilon_0} \left[1 - \left(1 - \frac{\epsilon_0}{\epsilon} \right) \frac{z}{l_z} \right]$

よって静電容量は, $C = Q/\Delta\phi = C_0 \frac{1}{1 - \left(1 - \frac{\epsilon_0}{\epsilon} \right) \frac{z}{l_z}}$

静電場のエネルギー

$U = \frac{1}{2} \epsilon_0 E_1^2 l_x l_y (l_z - z) + \frac{1}{2} \epsilon E_2^2 l_x l_y z$

$= \frac{1}{2} \epsilon_0 E_0^2 l_x l_y l_z \left[1 - \frac{z}{l_z} + \frac{\epsilon}{\epsilon_0} \left(\frac{E_2}{E_1} \right)^2 \frac{z}{l_z} \right]$

$= U_0 \left[1 - \left(1 - \frac{\epsilon_0}{\epsilon} \right) \frac{z}{l_z} \right]$

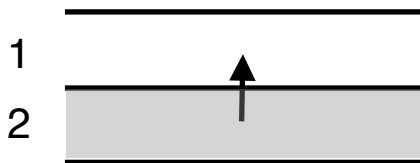
(U_0 は増減なし) $\frac{\epsilon_0}{\epsilon}$

よって境界面に働く力は

$F_z = - \frac{\partial U}{\partial z} = U_0 \left(1 - \frac{\epsilon_0}{\epsilon} \right) \frac{1}{l_z}$

単位面積あたり

$T = \frac{F_z}{l_x l_y} = \frac{U_0}{V} \left(1 - \frac{\epsilon_0}{\epsilon} \right)$



(3)

導体は等電位にある \rightarrow 領域 1, 2 の電場は同じ $E_1 = E_2 = E \hat{z}$

$$E_1 = \frac{\sigma_1}{\epsilon_0} \hat{z} \quad E_2 = \frac{\sigma_2}{\epsilon} \hat{z} \quad \text{なので} \quad \frac{\sigma_2}{\sigma_1} = \frac{\epsilon}{\epsilon_0} \dots \textcircled{1}$$

また $\sigma_1 \left(1 - \frac{x}{lx}\right) + \sigma_2 \frac{x}{lx} = \sigma_0$ (電荷の保存) $\dots \textcircled{2}$

電場は $E_1 = E_2 = E \hat{z}$

$$E = E_0 \frac{\sigma_1}{\sigma_0} = E_0 \frac{1}{1 + \left(\frac{\epsilon}{\epsilon_0} - 1\right) \frac{x}{lx}}$$

よから $\sigma_1 = \frac{\sigma_0}{1 + \left(\frac{\epsilon}{\epsilon_0} - 1\right) \frac{x}{lx}}$

電位差は

$$\Delta\phi = E l_2 = \frac{\sigma_1}{\epsilon_0} l_2$$

静電容量は

$$C = \frac{Q}{\Delta\phi} = C_0 \left[1 + \left(\frac{\epsilon}{\epsilon_0} - 1\right) \frac{x}{lx} \right]$$

$$= \frac{\sigma_0}{\epsilon_0} l_2 \frac{1}{1 + \left(\frac{\epsilon}{\epsilon_0} - 1\right) \frac{x}{lx}}$$

静電エネルギーは

$$U = \frac{1}{2} \epsilon E^2 x l_x l_y l_z + \frac{1}{2} \epsilon_0 E^2 (lx - x) l_y l_z$$

$$= \frac{1}{2} \epsilon_0 E_0^2 l_x l_y l_z \left(\frac{E}{E_0}\right)^2 \left[\frac{\epsilon}{\epsilon_0} \frac{x}{lx} + \left(1 - \frac{x}{lx}\right)\right]$$

U_0

$$= \frac{U_0}{1 + \left(\frac{\epsilon}{\epsilon_0} - 1\right) \frac{x}{lx}}$$

$$\leftarrow \frac{E}{E_0} = \frac{\sigma_1}{\sigma_0} = \left[1 + \left(\frac{\epsilon}{\epsilon_0} - 1\right) \frac{x}{lx} \right]^{-1}$$

境界面に働く力は

$$F_x = -\frac{\partial u}{\partial x} = \frac{U_0}{\left[1 + \left(\frac{\epsilon}{\epsilon_0} - 1\right) \frac{x}{l_x}\right]^2} \left(\frac{\epsilon}{\epsilon_0} - 1\right) \frac{1}{l_x}$$

単位面積あたりの大きさは

$$T = \frac{|F_x|}{l_y l_z} = \frac{U_0}{V_0} \frac{\frac{\epsilon}{\epsilon_0} - 1}{\left[1 + \left(\frac{\epsilon}{\epsilon_0} - 1\right) \frac{x}{l_x}\right]^2}$$

向きは図の向き。
($\epsilon > \epsilon_0$ のときは
xを増加する方向)

